

Funkce komplexní proměnné

Reziduová věta II

Použitím reziduové věty a pravidel pro počítání s rezidui spočtěte následující integrály: (ve smyslu Lebesgueově, Newtonově nebo ve smyslu hlavní hodnoty)

1. $\int_0^{2\pi} \frac{dt}{a + \cos t}, \quad a > 1$
2. $\int_0^{2\pi} \frac{dt}{(a + b \cos t)^2}, \quad a > b > 0$
3. $\int_0^{2\pi} \frac{\cos^2 t \, dt}{13 + 12 \cos t}$
4. $\int_0^\pi \frac{\cos^4 t \, dt}{1 + \sin^2 t}$
5. $\int_0^{2\pi} e^{\cos t} \cos((nt) - \sin t) \, dt, \quad n \in \mathbb{N}$
6. $\int_{-\pi}^\pi \frac{\sin(nt) \, dt}{1 - 2a \sin t + a^2}, \quad -1 < a < 1, n \in \mathbb{N}$
7. $\int_0^\infty \frac{dx}{(x+1)\sqrt{x}}$
8. $\int_0^\infty \frac{dx}{(x-1)\sqrt{x}}$
9. $\int_0^\infty \frac{dx}{x^a(x+b)}, \quad 0 < a < 1, b \neq 0$
10. $\int_0^\infty \frac{x^{a-1} dx}{(x+1)(x+2)(x+3)}, \quad 0 < a < 3$
11. $\int_0^1 \frac{x^{1-p}(1-x)^p dx}{(1+x)^2}, \quad -1 < p < 2$

$$12. \int_{-1}^1 \frac{(1+x)^{1-p}(1-x)^p dx}{x^2+1}, \quad -1 < p < 2$$

$$13. \int_0^1 \frac{x^{1-p}(1-x)^p dx}{x^2+1}, \quad -1 < p < 2$$

$$14. \int_{-1}^1 \ln \left(\frac{1+x}{1-x} \right) \frac{dx}{\sqrt[3]{(1-x)^2(1+x)}}$$

$$15. \int_0^\infty \frac{\ln x dx}{x^2+a^2}$$

$$16. \int_0^\infty \frac{\ln x dx}{(x-1)\sqrt{x}}$$

$$17. \int_0^\infty \frac{\ln x dx}{(x+1)^2 \sqrt[3]{x}}$$

Funkce komplexní proměnné

Reziduová věta III

Použitím reziduové věty a pravidel pro počítání s rezidui spočtěte následující integrály: (ve smyslu Lebesgueově, Newtonově nebo ve smyslu hlavní hodnoty)

$$1. \int_0^{\infty} \frac{\sin^2 x \, dx}{x^2}$$

$$2. \int_0^{\infty} \frac{\sin^3 x \, dx}{x^3}$$

$$3. \int_{-\infty}^{\infty} \frac{e^{ax} \, dx}{(e^x + 1)(e^x + 2)}, \quad 0 < a < 2$$

$$4. \int_0^{\infty} \frac{\sin(ax)}{\sinh x} \, dx, \quad a \in \mathbb{R}$$

$$5. \int_0^{\infty} \frac{\cosh(ax)}{\cosh(\pi x)} \, dx, \quad |a| < \pi$$

$$6. \int_0^{\infty} \frac{\sin^2(ax)}{\cosh x} \, dx, \quad a > 0$$

Výsledky následujících příkladů vyjádřete pomocí funkce Γ

$$7. \int_0^{\infty} x^z e^{-x^2} \, dx, \quad \operatorname{Re} z > -1$$

$$8. \int_0^1 \left(\ln \frac{1}{t}\right)^{z-1} dt, \quad \operatorname{Re} z > 0$$

$$9. \int_0^{\infty} t^{z-1} \cos t \, dt, \quad 0 < \operatorname{Re} z < 1$$

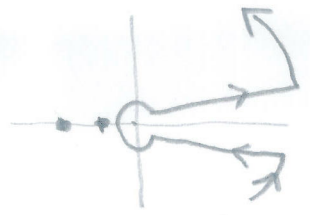
$$10. \int_0^{\infty} \cos x^p \, dx, \quad p > 1$$

$$11. \int_0^{\infty} \frac{\sin x^p}{x^p} \, dx, \quad p > \frac{1}{2}$$

11) $\int_0^1 \frac{x^{1-p} (1-x)^p}{(1+x)^2} dx, -1 < p < 2$

Uvedomme si: rozsah parametru p je dan
 požadavkem na konvergenci integrálu: $1-p > -1 \Rightarrow p < 2$
 a $p > -1$

Substitucí převedeme na známé případy: $\frac{x}{1-x} = y \Rightarrow x = \frac{y}{1+y}, 1+x = \frac{1+2y}{1+y}$
 $dy = \left(\frac{1}{1-x} + \frac{x^2}{(1-x)^2} \right) dx = \frac{1}{(1-x)^2} dx, 1-x = \frac{1}{1+y}$



$\Rightarrow I = \int_0^\infty y^{1-p} \cdot \frac{1}{1+y} \cdot \frac{1}{(1+2y)^2} \cdot \frac{1}{(1+y)^2} dy = \int_0^\infty \frac{y^{1-p}}{(1+y)(1+2y)^2} dy$. To už známé. Polky: $y = -1 \pm ix$

Parametrizace cesty opět $\gamma_1: z = t e^{i\theta} \quad t \in [0, R]$ $\gamma_3: z = t \cdot e^{i(2\pi-\theta)} \quad t \in [R, 0]$
 $\gamma_2: z = R e^{it} \quad t \in [0, 2\pi-\theta]$ $\gamma_4: z = \epsilon e^{it} \quad t \in [2\pi-\theta, \theta]$

$\gamma_2: M_R \sim \frac{R^{1-p}}{R^3} = R^{-2-p}$, tj. $M_R \sim R^{-1-p}$ a potřebujeme $-1-p < 0 \Rightarrow p > -1$

Jordanova lemma: $\int_{\gamma_2} \rightarrow 0$

γ_4 : dosadíme $z = \epsilon e^{it}$, použijeme definici integrálu a v mezinárodních ϵ dostaneme: $\epsilon \cdot \frac{\epsilon^{1-p}}{\epsilon^3} = \epsilon^{-1-p}$
 f je omezená na malé počátku \Rightarrow potřebujeme $\epsilon^{2-p} \rightarrow 0 \Rightarrow 2-p > 0 \Rightarrow p < 2$

$\int_{\gamma_1} \rightarrow I, \int_{\gamma_3} \frac{z^{1-p}}{(1+z)(1+2z)^2} dz \xrightarrow{z \rightarrow 0} \int_R^\epsilon \frac{t^{1-p} \cdot e^{2\pi i(1-p)}}{(1+t)(1+2t)^2} dt \xrightarrow{R \rightarrow \infty} -e^{-2\pi i(1-p)} I = -e^{-2\pi i p} I$

$\Rightarrow I(1 - e^{-2\pi i(1-p)}) = 2\pi i (\text{Res}_{-1} f + \text{Res}_{-1/2} f)$
 $\text{Res}_{-1} f = \frac{(-1)^{1-p}}{(-1)^2} = e^{\pi i(-1-p)} = -e^{-\pi i p}$
 $\text{Res}_{-1/2} f = \left(\frac{z^{1-p}}{(1+z)^2} \right)' \Big|_{z=-1/2} = \left[\frac{(1-p)z^{-p}}{(1+z)^2} - \frac{z^{1-p}}{(1+z)^3} \right]_{z=-1/2}$
 $= \frac{1-p}{2} \cdot \left(\frac{-1}{2}\right)^{-p} - \left(\frac{-1}{2}\right)^{(1-p)} = e^{-\pi i p} \left[\frac{1-p}{2^{1+p}} + \frac{1}{2^{1+p}} \right]$
 $\Rightarrow I = (-1)^{p-1} (2^p - 2^{p-1}) \cdot \frac{2\pi i}{e^{\pi i p} - e^{-\pi i p}} = \frac{\pi(2^p - p 2^{p-1} - 1)}{\sin \pi p}$

$p=0$ a $p=1$ třeba dopočítat zvlášť!

12) $I = \int_{-1}^1 \frac{(1+x)^{1-p} (1-x)^p}{x^2+1} dx, -1 < p < 2$. Stejně jako 11), $y = \frac{1+x}{1-x} \Rightarrow x = \frac{y-1}{y+1}$ $x \rightarrow -1 \Rightarrow y \rightarrow 0$
 $x \rightarrow 1 \Rightarrow y \rightarrow \infty$

$I = \int_0^\infty y^{1-p} \cdot \frac{2}{y+1} \cdot \frac{1}{(4-y^2+y^2)^2} \cdot \frac{2}{(y+1)^2} dy = \int_0^\infty \frac{y^{1-p}}{(y+1)(y^2+1)^2} dy$. Polky: $y = -1 \pm ix$

Cesta blížná, $\int_{\gamma_2} \rightarrow 0, \int_{\gamma_4} \rightarrow 0$

$I(1 - e^{-2\pi i p}) = 2\pi i (\text{Res}_{-1} f + \text{Res}_{-i} f + \text{Res}_{i} f) = 2\pi i \left(\frac{-2 \cdot e^{-\pi i p}}{2} + \frac{2i \cdot e^{-\pi i p}}{(1+i) \cdot 2i} + \frac{-2i \cdot e^{-\pi i p}}{(1-i) \cdot (-2i)} \right) = 2\pi i e^{-\pi i p} \left(\dots \right)$

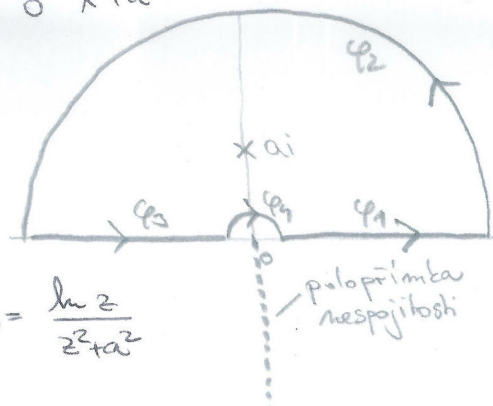
Platí: $\frac{e^{\pi i p}}{1+i} + \frac{e^{-\pi i p}}{1-i} = \frac{(e^{\pi i p} - \frac{-\pi i p}{2}) + i(e^{\frac{\pi i p}{2}} - e^{-\frac{\pi i p}{2}})}{2} = \cos \frac{\pi p}{2} + \sin \frac{\pi p}{2}$
 $\Rightarrow I = \frac{\pi}{\sin(\pi p)} \left[-1 + \cos \frac{\pi p}{2} + \sin \frac{\pi p}{2} \right]$

Opět $p=0$ a $p=1$ třeba dopočítat zvlášť!

13, D5

14, Uděláme níže, nejdříve 15, -17,

15) $\int_0^{\infty} \frac{\ln x dx}{x^2+a^2} = I$. Potvrďujeme, at' fce, která násobí $\ln x$ je sudá. Pak volíme tuto



Volíme komplexní logaritmus $\ln z = \ln|z| + i \arg z$, kde $\arg z \in (-\frac{\pi}{2}, \frac{3\pi}{2})$

$\gamma_1: z = t, t \in [-R, -\epsilon]$ $\gamma_3: z = t, t \in [\epsilon, R]$
 $\gamma_2: z = R e^{it}, t \in [0, \pi]$ $\gamma_4: z = \epsilon e^{it}, t \in [\pi, 0]$

$f(z) = \frac{\ln z}{z^2+a^2}$

$\gamma_2: M_R \sim \frac{1}{R^2} \ln R, RM_R \sim \frac{\ln R}{R}$ a dle Jordana

$\int_{\gamma_2} \rightarrow 0$

$\gamma_4: \text{Opit z definice integrálu dostáváme } \epsilon \cdot \ln \epsilon \cdot \max \frac{1}{z^2+a^2} \leq C$
 a protože $\epsilon \ln \epsilon \rightarrow 0, \int \rightarrow 0$

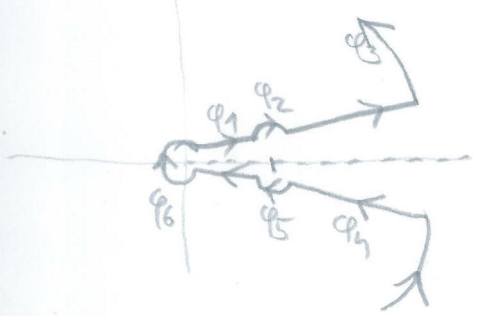
$\gamma_1: \int \rightarrow I$

$\gamma_3: \int \frac{\ln z}{z^2+a^2} = \int_{-R}^{-\epsilon} \frac{\ln t}{t^2+a^2} dt = \int_{-R}^{-\epsilon} \frac{\ln|t| + i\pi}{t^2+a^2} dt = \int_{\epsilon}^R \frac{\ln t}{t^2+a^2} dt + i\pi \int_{\epsilon}^R \frac{dt}{t^2+a^2}$

Celkem: $2I + \frac{i\pi}{2} \int_{-\infty}^{\infty} \frac{dt}{t^2+a^2} = 2\pi i \text{Res}_{ai} f$
 $2I + \frac{i\pi}{2} \cdot 2\pi i \text{Res}_{ai} \frac{1}{z^2+a^2} = 2\pi i \text{Res}_{ai} \frac{\ln z}{z^2+a^2}$

$I = \pi i \frac{\ln ai}{2ai} + \frac{\pi^2}{2} \cdot \frac{1}{2ai} = \frac{\pi}{2a} \left(\ln a + \frac{i\pi}{2} - \frac{i\pi}{2} \right) = \frac{\pi \ln a}{2a}$

16, $I = \int_0^{\infty} \frac{\ln x dx}{(x-1)\sqrt{x}}$. Fce bez $\ln x$ není sudá \Rightarrow postup z 15, nefunguje, Pomůže nám odmocnina. Polopřímka nespojitosti $\ln z$ bude zde kladná reálná poloosa.



$\int_{\gamma_1} f(z) dz \rightarrow I$

\int_{γ_2} : Očekáváme shora: LOOPN1: $\rightarrow iAb$, kde $b \rightarrow \pi$

$\gamma_2: A = \lim_{\text{shora}} \frac{\ln z}{\sqrt{z}} = \ln 1 = 0 \Rightarrow \int \rightarrow 0$

γ_3 : Jordan: $M_R \sim \frac{\ln R}{R^{3/2}}, RM_R \sim \frac{\ln R}{R^{1/2}} \rightarrow 0 \Rightarrow \int \rightarrow 0$

\int_{γ_4} : LOOPN2: $\rightarrow iAb$, kde $b \rightarrow -\pi$ a $A = \lim_{\text{zdola}} \frac{\ln z}{\sqrt{z}} = \frac{\ln 1 + 2\pi i}{-1} = -2\pi i$

$\Rightarrow \int_{\gamma_4} f \rightarrow i(-2\pi i)(-\pi) = -2\pi^2$

\int_{γ_6} : Definice křivkového int: $\int_{\gamma_6} |f| \leq \frac{\epsilon \ln \epsilon}{\sqrt{\epsilon}} \cdot \max \left| \frac{1}{z-1} \right| \xrightarrow{\epsilon \rightarrow 0} 0$

$f \rightarrow \ln z \rightarrow \ln t + 2\pi i$
 $\sqrt{z} \rightarrow \sqrt{t} \cdot e^{\frac{2\pi i}{2}} = -\sqrt{t}$

$\int_{\gamma_4} f \rightarrow \int_{\gamma_4} \frac{\ln t + 2\pi i}{(t-1) \cdot (-\sqrt{t})} dt = I + 2\pi i \int_0^{\infty} \frac{dt}{\sqrt{t}(t-1)}$

opacný smer

Celkem: $2I + 2\pi i \int_0^{\infty} \frac{dt}{\sqrt{t}(t-1)} - 2\pi^2 = 0$ (f je holomorfní vnitřně Γ)
 ale pozor ve smyslu hlavních hodnot!!

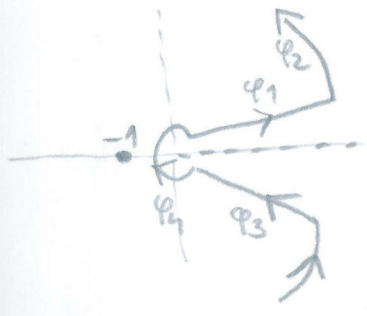
Zbývá integrál šed' opět reziduovou větou nebo přímo: $y = \sqrt{t}, dy = \frac{1}{2\sqrt{t}} dt$

$\Rightarrow J = \int_0^{\infty} \frac{2dy}{y^2-1} = \int_0^{\infty} \left(\frac{1}{y-1} - \frac{1}{y+1} \right) dy$

$\lim_{\epsilon \rightarrow 0} \left[\ln \frac{1-y}{1+y} \right]_0^{1-\epsilon} + \left[\ln \frac{y-1}{y+1} \right]_{1+\epsilon}^{\infty} = \lim_{\epsilon \rightarrow 0^+} \frac{\ln \epsilon}{2-\epsilon} - \frac{\ln \epsilon}{2+\epsilon}$

$= \lim_{\epsilon \rightarrow 0^+} \frac{2\epsilon \ln \epsilon}{4-\epsilon^2} = 0 \Rightarrow J = 0$ a $I = \pi^2$

17) $I = \int_0^{\infty} \frac{\ln x dx}{(x+1)^2 \sqrt[3]{x}}$



$f(z) = \frac{\ln z}{(z+1)^2 z^{1/3}}$ Nespojnost na \mathbb{R}^+

$\int_{\gamma_2} : \text{Jordan } M_{\epsilon} \sim \frac{\ln R}{R^{2+1/3}}, RM_{\epsilon} \sim \frac{\ln R}{R^{4/3}} \rightarrow 0 \Rightarrow \int_{\gamma_2} \rightarrow 0$

$\int_{\gamma_4} : \text{Máme } \epsilon \cdot \ln \epsilon \cdot \epsilon^{-1/3} \cdot \max_{|z+1| \leq \epsilon} \left| \frac{1}{(z+1)^2} \right| \xrightarrow{\epsilon \rightarrow 0} 0 \Rightarrow \int_{\gamma_4} \rightarrow 0$

$\int_{\gamma_3} : \ln z \rightarrow \ln t + 2\pi i$
 $\sqrt[3]{z} \rightarrow \sqrt[3]{t} \cdot e^{\frac{2\pi i}{3}}$

$\int_{\gamma_3} f \rightarrow - \int_0^{\infty} \frac{\ln t + 2\pi i}{(t+1)^2 \sqrt[3]{t}} \cdot e^{-\frac{2\pi i}{3}} dt$

$= -e^{-\frac{2\pi i}{3}} I - 2\pi i e^{-\frac{2\pi i}{3}} \int_0^{\infty} \frac{1}{(t+1)^2 \sqrt[3]{t}} dt$

$\Rightarrow I \cdot (1 - e^{-\frac{2\pi i}{3}}) - 2\pi i e^{-\frac{2\pi i}{3}} \int_0^{\infty} \frac{dt}{(t+1)^2 \sqrt[3]{t}} = 2\pi i \text{Res}_{-1} f = 2\pi i \left(\frac{\ln z}{\sqrt[3]{z}} \right)' \Big|_{z=-1}$

$= 2\pi i \left[\frac{1}{z^3 \sqrt[3]{z}} - \frac{1}{3} \frac{\ln z}{z^{4/3}} \right] \Big|_{z=-1} = 2\pi i (-1) \cdot e^{\frac{\pi i}{3}} \cdot \left(1 - \frac{i\pi}{3} \right)$

$= -2\pi i e^{-\frac{\pi i}{3}} \cdot \left(1 - \frac{i\pi}{3} \right)$

J... opět reziduovou větou, stejná cesta

$J \cdot (1 - e^{-\frac{2\pi i}{3}}) = 2\pi i \cdot \left(\frac{z^{-1/3}}{z^3} \right)' \Big|_{z=-1} = 2\pi i \cdot \left(-\frac{1}{3} \right) \cdot (-1)^{-4/3} = \frac{2\pi i}{3} \cdot e^{-\frac{\pi i}{3}}$

$\Rightarrow J = \frac{\pi}{3 \cdot \sin \frac{\pi}{3}} = \frac{2\pi}{3\sqrt{3}}$

$\Rightarrow I \cdot \left(\sin \frac{\pi}{3} \right) - \pi \cdot e^{-\frac{\pi i}{3}} \cdot \frac{2\pi}{3\sqrt{3}} = -\pi \cdot \left(1 - \frac{i\pi}{3} \right)$

$I \cdot \frac{\sqrt{3}}{2} - \frac{2\pi^2}{3\sqrt{3}} \cdot \left(\frac{1}{2} - \frac{\sqrt{3}}{2}i \right) = -\pi + \frac{i\pi^2}{3} \Rightarrow \underline{\underline{I = -\frac{2\pi}{\sqrt{3}} + \frac{2\pi^2}{9}}}$

$$14) I = \int_{-1}^1 \ln\left(\frac{1+x}{1-x}\right) \frac{dx}{\sqrt{(1-x)^2(1+x)}}$$

Substituce $y = \frac{1+x}{1-x}$ $x \rightarrow -1 \Rightarrow y \rightarrow 0$
 $x \rightarrow 1 \Rightarrow y \rightarrow +\infty$

$$dy = \left(\frac{1}{1-x} + \frac{1+x}{(1-x)^2}\right) dx = \frac{2}{(1-x)^2} dx$$

$$x = \frac{y-1}{y+1} \quad 1-x = \frac{2}{y+1}$$

$$I = \int_0^\infty \ln y \cdot y^{1/3} \cdot \left(\frac{y+1}{2}\right) \cdot \frac{4}{(y+1)^2} \cdot \frac{1}{2} dy$$

$$= \int_0^\infty \frac{\ln y \, dy}{y^{1/3} \cdot (y+1)}$$

To je také jako příklad 17, jen $\nu = y = -1$ je jednoduchý pól místo dvojnásobného. Stejná cesta, stejný postup. Dostaneme

$$I \cdot (1 - e^{-2\pi i/3}) - 2\pi i e^{-2\pi i/3} \int_0^\infty \frac{dt}{t^{1/3}(t+1)} = 2\pi i \operatorname{Res}_{-1} f = 2\pi i \frac{\ln(-1)}{(-1)^{1/3}} = 2\pi i \left(\frac{i\pi \cdot e^{-i\pi/3}}{e^{-i\pi/3}} \right) = -2\pi^2 \cdot e^{-i\pi/3}$$

$$J = \int_0^\infty \frac{dt}{t^{1/3}(t+1)}$$

buď spočítáme stejnou integrační cestou nebo rovnou použijeme výsledek příkladu 9)
 $s = a = 1/3, b = 1 \Rightarrow J = \frac{\pi}{\sin \pi/3} = \frac{2\pi}{\sqrt{3}}$

$$\Rightarrow I \cdot 2i \sin \frac{\pi}{3} - 2\pi i \cdot \frac{2\pi}{\sqrt{3}} \cdot \left(\cos \frac{\pi}{3} - i \sin \frac{\pi}{3} \right) = -2\pi^2$$

$$I \cdot 2i \cdot \frac{\sqrt{3}}{2} = i \frac{2\pi^2}{\sqrt{3}} \Rightarrow \underline{\underline{I = \frac{2}{3}\pi^2}}$$

Poslední sada příkladů: různé méně standardní příklady, kde se naučíme další triky

$$1) I = \int_0^\infty \frac{\sin^2 x}{x^2} dx$$

... suda'

$$\sin x = \frac{e^{ix} - e^{-ix}}{2i} \Rightarrow \sin^2 x = \frac{e^{2ix} - 2 + e^{-2ix}}{-4}$$

Nelze počítat na jedné křivce Jordanovi, viz níže

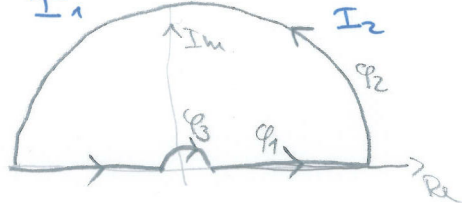
$$\Rightarrow I = -\frac{1}{8} \left(\int_{-\infty}^\infty \frac{e^{2ix} - 1}{x^2} dx + \int_{-\infty}^\infty \frac{e^{-2ix} - 1}{x^2} dx \right)$$

Přivedení fce $\frac{\sin^2 x}{x^2}$ má $\nu = x=0$ konečnou limitu

Nové fce $\frac{e^{\pm 2ix} - 1}{x^2}$ mají $\nu = x=0$ pól, ale nastěší jen 1-násobný, protože

$$\lim_{x \rightarrow 0} x \cdot \frac{e^{\pm 2ix} - 1}{x^2} = \pm 2i \neq 0$$

$$I_1: f(z) = \frac{e^{2iz} - 1}{z^2}$$



$$\lim_{R \rightarrow \infty} \int_{\Gamma_1} f(z) dz = \text{p.v.} \int_{-\infty}^\infty f(z) dz = I_1$$

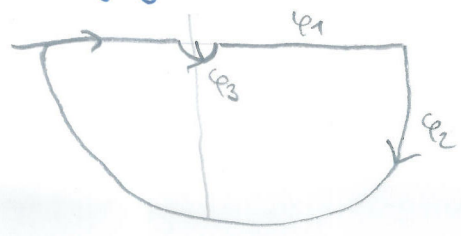
Γ_2 : Jordanova lemma $\left\{ \begin{array}{l} \text{pro fci } -\frac{1}{z^2} \text{ je } \alpha=0, M_R \sim \frac{1}{R^2}, R \rightarrow \infty \rightarrow 0 \text{ OK} \\ \text{pro fci } \frac{e^{2iz}}{z^2} \text{ je } \alpha=2, M_R \sim \frac{1}{R^2} \rightarrow 0 \text{ OK} \end{array} \right\} \int_{\Gamma_2} f \rightarrow 0$

$$\Gamma_3: \int_{\Gamma_3} f(z) dz \xrightarrow{\epsilon \rightarrow 0} iAb \text{ dle LOTN1, kde } b = -\pi \text{ a } A = \lim_{z \rightarrow 0} z \cdot \frac{e^{2iz} - 1}{z^2} = 2i$$

$$\Rightarrow \int_{\Gamma_3} f(z) dz \rightarrow i(2i) \cdot (-\pi) = 2\pi \Rightarrow I_1 + 0 + 2\pi = 0 \Rightarrow \underline{\underline{I_1 = -2\pi}}$$

Candijova věta, f je holomorfní

I_2 analogicky, jen pro Jordana potřeby je $\text{Im } z < 0$, potřebujeme e^{-2iz}



Opět $\int_{\gamma_1} f(z) dz \rightarrow I_2$

$\int_{\gamma_2} f$ dle Jordana $\left\{ \begin{array}{l} \alpha=0, M_R \sim \frac{1}{R^2} \text{ pro } -\frac{1}{z^2} \\ \alpha=-2, M_R \sim \frac{1}{R^2} \text{ pro } \frac{e^{-2iz}}{z^2} \end{array} \right\} \rightarrow 0$

$I_2 + 0 + 2\pi = 0$ ← Cauchy

$\int_{\gamma_3} f \rightarrow iAb$ dle Cauchy, kde $b = \pi$, $A = \lim_{z \rightarrow 0} z \cdot \frac{e^{-2iz} - 1}{z^2} = -2i$

$\Rightarrow I_2 = -2\pi$

$\Rightarrow \int_{\gamma_3} f \rightarrow i \cdot (-2i) \cdot \pi = 2\pi$

$\Rightarrow I = -\frac{1}{8} (-2\pi - 2\pi) = \frac{\pi}{2}$

2) $\int_0^{\infty} \frac{\sin^3 x}{x^3} = I$

$\sin^3 x = -\frac{1}{8i} (e^{3ix} - 3e^{ix} + 3e^{-ix} - e^{-3ix})$

$= -\frac{1}{8i} (e^{ix} \cdot (e^{2ix} - 1) - 2e^{ix} + e^{-ix} \cdot (e^{-2ix} - 1) + 2e^{-ix})$

$\frac{e^{2ix} - 1}{x^3}$ má pořadí pólu má s. 2 \Rightarrow musíme upravit dál:

$= -\frac{1}{8i} (e^{ix} (e^{2ix} - 1 - 2ix) + 2e^{ix} (e^{-ix} + ix - 1) - e^{-ix} (e^{-2ix} - 1 + 2ix) + 2e^{-ix} (e^{ix} - 1 - ix))$

Tedy jsem vytvořil póly násobnosti 1.

Opět ze sudosti $I = \frac{1}{2} \int_{-\infty}^{\infty} \frac{\sin^3 x}{x^3} dx = -\frac{1}{16i} (I_1 + I_2 + I_3 + I_4)$

$I_1 = \int_{-\infty}^{\infty} e^{ix} \cdot \frac{e^{2ix} - 1 - 2ix}{x^3} dx$... integrujeme přes horní půlkruh s obcházením pólu π nule

$\int_{\gamma_1} f \rightarrow I_1$, \int_{γ_2} Jordan: $\frac{e^{3iz}}{z^3}$ OK, $\alpha=3, M_R \sim \frac{1}{R^3}$, $-\frac{e^{iz}}{z^3}$ OK, $\alpha=1, M_R \sim \frac{1}{R^3}$, $-\frac{2ie^{iz}}{z^2}$ OK, $\alpha=1, M_R \sim \frac{1}{R^2}$

$\int_{\gamma_3} f \rightarrow iAb$ dle Cauchy: $b = -\pi$, $A = \lim_{z \rightarrow 0} \underbrace{e^{iz}}_1 \cdot \underbrace{\frac{e^{2iz} - 1 - 2iz}{z^3}}_{\text{Taylor}} \cdot z = \frac{(2i)^2}{2!} = -2$

$\Rightarrow I_1 + 0 + 2\pi i = 0$ dle Cauchy (f holomorfní) $\Rightarrow I_1 = -2\pi i$

$I_2 = \int_{-\infty}^{\infty} \frac{2e^{ix} (e^{-ix} + ix - 1)}{x^3} dx$... Stejná cesta: $\int_{\gamma_1} f \rightarrow I_2$

$\int_{\gamma_2} f$ Jordan: $\frac{2e^{iz}}{z^3}$ OK, $\alpha=0, M_R \sim \frac{1}{R^3}$, $\frac{2ie^{iz}}{z^2}$ OK, $\alpha=1, M_R \sim \frac{1}{R^2}$, $-\frac{2e^{iz}}{z^3}$ OK, $\alpha=1, M_R \sim \frac{1}{R^3}$

$\Rightarrow \int_{\gamma_2} f \rightarrow 0$

$\int_{\gamma_3} f \rightarrow iAb$ dle Cauchy: $b = -\pi$, $A = \lim_{z \rightarrow 0} 2 \cdot e^{iz} \cdot \frac{e^{-iz} - 1 + iz}{z^3} \cdot z = 2 \cdot 1 \cdot \frac{(-i)^2}{2!} = -1$

$\Rightarrow I_2 + 0 + \pi i = 0 \Rightarrow I_2 = -\pi i$

I_3 a I_4 analogicky, jen musíme kvůli Jordanovi integrovat přes dolní polovosinu

Podobně jako v předchozím příkladě.

$$I_3: \int_{\gamma_1} \rightarrow I_3, \int_{\gamma_2} \text{Jordan} \rightarrow 0, \int_{\gamma_3} f \rightarrow iAb, b=\pi, A = \lim_{z \rightarrow 0} -e^{-iz} \cdot \left(\frac{e^{-2iz} - 1 + 2iz}{z^3} \right) \cdot z =$$

$$= - \frac{(-2i)^2}{2!} = 2$$

$$\Rightarrow I_3 + 0 + 2\pi i = 0 \Rightarrow I_3 = -2\pi i$$

$$I_4: \int_{\gamma_1} \rightarrow I_4, \int_{\gamma_2} \text{Jordan} \rightarrow 0, \int_{\gamma_3} f \rightarrow iAb, b=\pi, A = \lim_{z \rightarrow 0} -2 \cdot e^{-iz} \cdot \left(\frac{e^{iz} - 1 - iz}{z^3} \right) \cdot z =$$

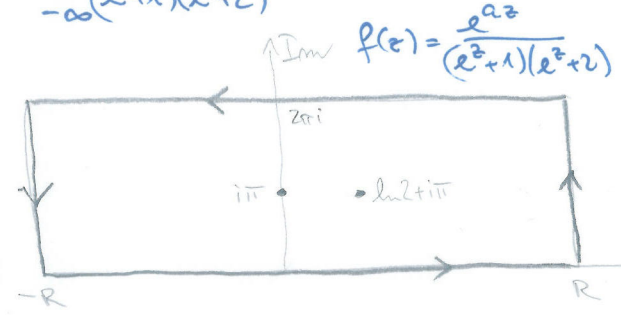
$$= -2 \cdot \frac{1^2}{2!} = 1$$

$$\Rightarrow I_4 + 0 + \pi i = 0 \Rightarrow I_4 = -\pi i$$

$$\text{Celkem } I = -\frac{1}{16i} (-6\pi i) = \underline{\underline{\frac{3}{8}\pi}}$$

$$3) \int_{-\infty}^{\infty} \frac{e^{ax}}{(x^2+1)(x^2+2)} dx = I \quad 0 < a < 2$$

Póly: $e^x = -1 = e^{i\pi + 2k\pi i} \Rightarrow x = i(\pi + 2k\pi)$
 $e^x = -2 = e^{\ln 2 + i\pi + 2k\pi i} \Rightarrow x = \ln 2 + i(\pi + 2k\pi)$



$$\begin{aligned} \gamma_1(t) &= t, \quad t \in [-R, R] \\ \gamma_2(t) &= R + it, \quad t \in [0, 2\pi] \\ \gamma_3(t) &= t + 2\pi i, \quad t \in [R, -R] \\ \gamma_4(t) &= -R + it, \quad t \in [2\pi, 0] \end{aligned}$$

$$\int_{\gamma_1} f(z) dz \rightarrow I$$

$$\int_{\gamma_2} f(z) dz = \int_0^{2\pi} \frac{e^{a(R+it)}}{(e^{2R+2it}+1)(e^{2R+2it}+2)} \cdot i dt$$

$$|f| \leq \int_0^{2\pi} \frac{e^{aR}}{(e^{2R}-1)(e^{2R}-2)} dt \leq 2\pi \cdot e^{R(a-2)} \xrightarrow{R \rightarrow \infty} 0 \text{ protože } a < 2$$

$$\int_{\gamma_3} f(z) dz = - \int_{-R}^R \frac{e^{a(t+2\pi i)}}{(e^{2t+4\pi i}+1)(e^{2t+4\pi i}+2)} dt = -e^{2\pi ia} \int_{\gamma_1} f(z) dz \rightarrow -e^{2\pi ia} I$$

$$\int_{\gamma_4} f(z) dz = - \int_0^{2\pi} \frac{e^{a(-R+it)}}{(e^{-2R+2it}+1)(e^{-2R+2it}+2)} \cdot i dt \Rightarrow |f| \leq \int_0^{2\pi} \frac{e^{-aR}}{(1-e^{-2R})(2-e^{-R})} dt \leq 4\pi e^{-aR} \rightarrow 0 \text{ protože } a > 0$$

$$\text{Celkem: } I \cdot (1 - e^{2\pi ia}) = (\text{Res}_{i\pi} f + \text{Res}_{\ln 2 + i\pi} f) \cdot 2\pi i$$

$$= \cancel{\dots} = \pi (1 - 2^{a-1})$$

Pozor na výpočet reziduí! Nemáme tvar $z - z_0$ ve jmenovateli!

$$\text{Res}_{i\pi} f = \frac{e^{ai\pi}}{(e^{i\pi}+2)} \cdot \frac{1}{(z^2+1)'} \Big|_{z=i\pi} = \frac{e^{ai\pi}}{1} \cdot \frac{1}{e^{i\pi}} = -e^{ai\pi}$$

$$\text{Podobně } \text{Res}_{\ln 2 + i\pi} f = \frac{e^{a(\ln 2 + i\pi)}}{-1} \cdot \frac{1}{(z^2+2)'} \Big|_{z=\ln 2 + i\pi} = -2 e^{a(\ln 2 + i\pi)} \cdot \left(-\frac{1}{2}\right) = 2^{a-1} e^{ai\pi}$$

$$I(1 - e^{2\pi ia}) = \pi (1 - 2^{a-1})$$

$$I = \frac{\pi}{\sin(\pi a)} \cdot (1 - 2^{a-1})$$

To all nefunguje pro $a=1$. Tam musíme počítat metodami 1. semestru

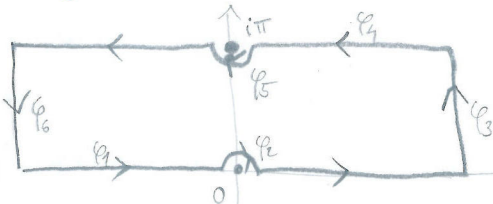
$$\int_{-\infty}^{\infty} \frac{e^x dx}{(e^x+1)(e^x+2)} \quad \left. \begin{array}{l} t=e^x \quad x \rightarrow -\infty \Rightarrow t \rightarrow 0 \\ dt=e^x dx \quad x \rightarrow +\infty \Rightarrow t \rightarrow \infty \end{array} \right\} \Rightarrow I = \int_0^{\infty} \frac{dt}{(t+1)(t+2)}$$

$$\frac{1}{(t+1)(t+2)} = \frac{A}{t+1} + \frac{B}{t+2} \Rightarrow A=1, B=-1$$

$$\Rightarrow I = \left[\ln \frac{t+1}{t+2} \right]_0^{\infty} = \ln 1 - \ln \frac{1}{2} = \ln 2$$

4) $I = \int_0^{\infty} \frac{\sin(ax)}{\sin bx} dx, a \in \mathbb{R}$. Sudost: $I = \frac{1}{2} \int_{-\infty}^{\infty} \frac{\sin(ax)}{\sin bx} dx = \frac{1}{2} \operatorname{Im} \int_{-\infty}^{\infty} \frac{e^{iax}}{e^{ix} - e^{-ix}} dx$
 $= \operatorname{Im} \int_{-\infty}^{\infty} \frac{e^{iax}}{e^{ix} - e^{-ix}} dx$... zde už integrál jako p.v.

Póly $f(z) = \frac{e^{iaz}}{e^z - e^{-z}}$: $e^z = e^{-z} \Leftrightarrow e^{2z} = 1 = e^{2k\pi i} \Rightarrow$ póly jsou $z_k = k\pi i, k \in \mathbb{Z}$



$\gamma_1: z=t, t \in [-R, -\epsilon] \cup [\epsilon, R]$
 $\gamma_2: z = \epsilon e^{it}, t \in [\pi, 0]$
 $\gamma_3: z = R + it, t \in [0, \pi]$
 $\gamma_4: z = t, t \in [R, \epsilon] \cup [-\epsilon, -R]$
 $\gamma_5: z = i\pi + \epsilon e^{it}, t \in [0, -\pi]$
 $\gamma_6: z = -R + it, t \in [\pi, 0]$

$\int_{\gamma_1} f \rightarrow \int_{\gamma_2} f \rightarrow iAb$ dle LOOPN1, $b = -\pi, A = \lim_{z \rightarrow 0} z f(z) = \lim_{z \rightarrow 0} \frac{z}{e^z - e^{-z}} = \lim_{z \rightarrow 0} \frac{1}{e^z} \cdot \frac{z}{e^z - 1} \cdot \frac{1}{z} = \frac{1}{2}$

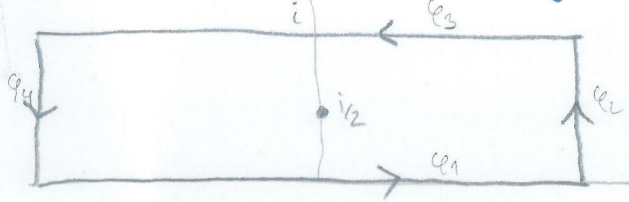
$\left| \int_{\gamma_3} f \right| \leq \int_0^{\pi} \frac{e^{-at}}{e^R \cdot |e^{it} - e^{-2R+it}|} dt \leq \frac{C}{e^R} \rightarrow 0$, $\int_{\gamma_4} f \rightarrow \int_{-\infty}^{\infty} \frac{e^{ia(t+i\pi)}}{e^{t+i\pi} - e^{-t-i\pi}} dt = e^{-a\pi}$

$\int_{\gamma_2} f \rightarrow iAb$ dle LOOPN1: $b = -\pi, A = \lim_{z \rightarrow i\pi} z f(z) = (z - i\pi) = \lim_{z \rightarrow i\pi} \frac{z - i\pi}{e^z - e^{-z}} \stackrel{\text{Hosp}}{=} e^{-a\pi} \cdot \lim_{z \rightarrow i\pi} \frac{1}{e^z + e^{-z}} = -\frac{e^{-a\pi}}{2}$

$\left| \int_{\gamma_6} f \right| \leq \int_0^{\pi} \frac{e^{-at}}{e^R \cdot |e^{it} - e^{-2R+it}|} dt \leq \frac{C}{e^R} \rightarrow 0$
 DOHROMADY: $\int (1 + e^{-a\pi}) - \frac{i\pi}{2} (1 - e^{-a\pi}) = 0$
 $\int = \frac{i\pi}{2} \left(\frac{1 - e^{-a\pi}}{1 + e^{-a\pi}} \right)$
 $I = \operatorname{Im} \int = \frac{\pi}{2} \cdot \tanh \frac{a\pi}{2}$

5) $I = \int_0^{\infty} \frac{\cosh(ax)}{\cosh(\pi x)} dx, |a| < \pi$. Sudost: $I = \frac{1}{2} \int_{-\infty}^{\infty} \frac{e^{ax} + e^{-ax}}{e^{\pi x} + e^{-\pi x}} dx$. Označme $J_a = \int_{-\infty}^{\infty} \frac{e^{ax}}{e^{\pi x} + e^{-\pi x}} dx$
 Lehe nahledneme, že $J_a = J_{-a}$: V J_a uděláme substituci $y = -x; dy = -dx, J_a = \int_{+\infty}^{-\infty} \frac{e^{-ay}}{e^{\pi y} + e^{-\pi y}} \cdot (-1) dy = \int_{-\infty}^{+\infty} \frac{e^{-ay}}{e^{\pi y} + e^{-\pi y}} dy = J_{-a}$
 Protože $I = \frac{1}{2}(J_a + J_{-a})$, je tak $I = J_a$. $f(z) := \frac{e^{az}}{e^{\pi z} + e^{-\pi z}}$

Póly f : $e^{\pi z} + e^{-\pi z} = 0 \Rightarrow e^{2\pi z} = -1 = e^{i\pi + 2k\pi i} \Rightarrow z_k = \frac{i}{2} + k i$



$\gamma_1: z=t, t \in [-R, R]$
 $\gamma_2: z = R + it, t \in [0, 1]$
 $\gamma_3: z = t, t \in [R, -R]$
 $\gamma_4: z = -R + it, t \in [1, 0]$

$$\int_{\varphi_1} f \rightarrow J_a, \quad \int_{\varphi_3} f(z) dz = \int_{-R}^R \frac{e^{a(i+t)}}{e^{\pi(i+t)} + e^{-\pi(i+t)}} dt = e^{ai} \int_{-R}^R \frac{e^{at}}{e^{t+\pi} + e^{-t-\pi}} dt \rightarrow e^{ai} J_a$$

$$|\int_{\varphi_2} f| \leq \int_0^1 \frac{e^{aR}}{e^{\pi R} \cdot |e^{\pi it} + e^{-2\pi R - \pi it}|} dt \leq C e^{(a-\pi)R} \rightarrow 0 \text{ pro } a < \pi$$

$$|\int_{\varphi_4} f| \leq \int_0^1 \frac{e^{-aR}}{e^{\pi R} \cdot |e^{-\pi it} + e^{-2\pi R - \pi it}|} dt \leq C \cdot e^{-\pi R} \rightarrow 0$$

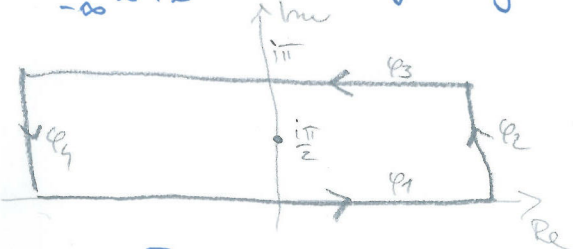
Celkem: $(1+e^{ai}) J_a = (\text{Res}_{i/2} f) \cdot 2\pi i = 2\pi i \frac{e^{ai/2}}{\pi \cdot e^{i\pi/2} - \pi \cdot e^{-i\pi/2}} = 2\pi i \frac{e^{ai/2}}{2\pi i} = e^{ai/2}$

$\Rightarrow J_a = \frac{1}{2 \cos \frac{a}{2}}$ pro $a < \pi$. Aby bylo J_a také dobře definováno, musí být i $-a < \pi \Rightarrow |a| < \pi$

6) $I = \int_0^{\infty} \frac{\sin^2 ax}{\cosh x} dx, a > 0$ Sudost: $I = \frac{1}{2} \int_{-\infty}^{\infty} \frac{\sin^2 ax}{\cosh x} dx = \frac{1}{2} \int_{-\infty}^{\infty} \frac{1 - \cos 2ax}{e^x + e^{-x}} dx$

$= \frac{1}{2} \int_{-\infty}^{\infty} \frac{1}{e^x + e^{-x}} dx - \frac{1}{2} \int_{-\infty}^{\infty} \text{Re} \frac{e^{2iax}}{e^x + e^{-x}} dx$ $f(z) := \frac{e^{2iaz}}{e^z + e^{-z}}$ Póly: jako v př. 5, $z_k = \frac{\pi i}{2} + k\pi i$

$J_a := \int_{-\infty}^{\infty} \frac{e^{2iax}}{e^x + e^{-x}} dx$ Integracní cesta podobně jako v př. 5,



$$\int_{\varphi_1} f \rightarrow J_a, \quad \int_{\varphi_3} f(z) dz = \int_{-R}^R \frac{e^{2ia(i\pi+t)}}{e^{i\pi \cdot t} + e^{-i\pi \cdot t}} dt = e^{-2a\pi} \int_{-R}^R \frac{e^{2iat}}{e^t + e^{-t}} dt \rightarrow e^{-2a\pi} J_a$$

$$|\int_{\varphi_2} f| \leq \int_0^{\pi} \frac{e^{-2at}}{e^R |e^{it} + e^{-2R - it}|} dt \leq C \cdot e^{-R} \rightarrow 0$$

$$|\int_{\varphi_4} f| \leq \int_0^{\pi} \frac{e^{-2at}}{e^R |e^{-it} + e^{-2R + it}|} dt \leq C \cdot e^{-R} \rightarrow 0$$

Celkem: $J_a (1 + e^{-2a\pi}) = 2\pi i \text{Res}_{i\pi/2} f$
 $= 2\pi i \frac{e^{-a\pi}}{e^{i\pi/2} - e^{-i\pi/2}} = \pi e^{-a\pi}$
 $\Rightarrow J_a = \frac{\pi}{2 \cosh a\pi}$

To platí: pro $a=0!$ $J_0 = \frac{1}{2}$

$$I = \frac{1}{2} (J_0 - \text{Re } J_a) = \frac{1}{2} \left(\frac{\pi}{2} - \frac{\pi}{2 \cosh a\pi} \right) = \frac{\pi}{4} \left(1 - \frac{1}{\cosh a\pi} \right)$$