

Maxwellove rovnice

$$\begin{aligned}\nabla \cdot \vec{E} &= \frac{\rho}{\epsilon_0} & \nabla \times \vec{E} &= -\frac{\partial \vec{B}}{\partial t} \\ \nabla \cdot \vec{B} &= 0 & \nabla \times \vec{B} &= \mu_0 \vec{j} + \frac{1}{c^2} \frac{\partial \vec{E}}{\partial t}\end{aligned}$$

Transformácie polí

$$F'_{\mu\nu} = \Lambda_{\mu}^{\mu'} F_{\mu'\nu'} \Lambda_{\nu}^{\nu'}$$

$$\begin{aligned}E'_{\parallel} &= E_{\parallel} & E'_{\perp} &= (E_{\perp} + \frac{v}{c} \times c\vec{B})\gamma \\ B'_{\parallel} &= B_{\parallel} & B'_{\perp} &= (cB_{\perp} - \frac{v}{c} \times \vec{E})\gamma\end{aligned}$$

$$\begin{aligned}\phi' &= \gamma \left(\phi - \frac{v}{c} A_{\parallel} \right) \\ A'_{\parallel} &= \gamma \left(A_{\parallel} - \frac{v}{c} \phi \right)\end{aligned}$$

Dipól

$$\vec{m} = \frac{1}{2} \int \vec{r} \times \vec{j} dV$$

$$\vec{A} = \frac{\mu_0}{4\pi} \frac{\vec{m} \times \vec{r}}{r^3}$$

$$\vec{F} = \vec{m} \cdot \nabla \vec{B}$$

$$\vec{M} = \vec{m} \times \vec{B}$$

$$b = \nabla \times a$$

$$b^1 = \frac{1}{h_1 h_2 h_3} \left((h_3^2 a^3)_{12} - (h_2^2 a^2)_{13} \right)$$

Biot-Savart

$$\vec{A} = \frac{\mu_0}{4\pi} \int \frac{\vec{j}'}{|\vec{x} - \vec{x}'|} dV$$

$$\vec{B} = \frac{\mu_0}{4\pi} \int \vec{j}'(\vec{x}') \times \frac{\vec{x} - \vec{x}'}{|\vec{x} - \vec{x}'|^3} dV$$

$$\vec{B} = \frac{\mu_0}{4\pi} I \int d\vec{s} \times \frac{\vec{x} - \vec{x}'}{|\vec{x} - \vec{x}'|^3}$$

Integrálny Ampér

$$\oint_{\Gamma} \vec{B} \cdot d\vec{l} = \mu_0 I$$

Magnetická sila

$$\vec{F} = \int \vec{j} \times \vec{B} dV$$

Indukčnosť

$$L = \frac{\mu_0}{4\pi} \oint_{\Gamma_k} \oint_{\Gamma_l} \frac{d\vec{s}_k \cdot d\vec{s}_l}{|\vec{r}_k - \vec{r}_l|}$$

$$U = \frac{1}{2} L I^2 \rightarrow 1 \text{ smyčka}$$

$$U = \frac{1}{2} I_k L_{kl} I_l \rightarrow \text{viacer smyčiek}$$

$$\Psi_k = L_{kl} I_l$$

Magnetický tok

$$\Psi = \int \vec{B} \cdot d\vec{S}$$

$$-\frac{d\Psi}{dt} = \mathcal{E}$$

Energia

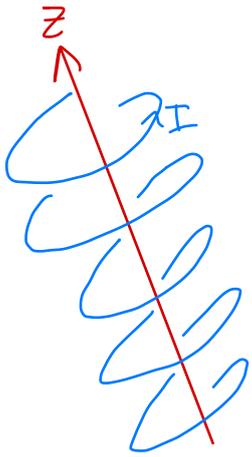
$$U_E = \frac{1}{2} \epsilon_0 E^2$$

$$U_B = \frac{1}{2} \epsilon_0 B^2$$

$$\vec{S} = \epsilon_0 c^2 \vec{E} \times \vec{B}$$

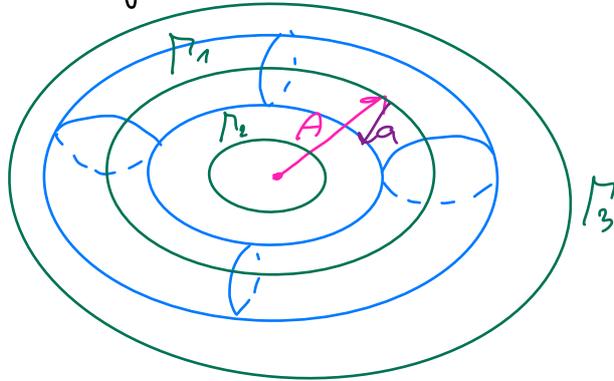
$$\vec{j} = \frac{1}{c^2} \vec{S}$$

10.2 Pole nekonečného solenoidu



$$n = \frac{N}{L} \dots \text{hust. vinutia}$$

$$\oint \vec{B} \cdot d\vec{\ell} = \mu_0 I$$



$$\int_{\Gamma_2} \vec{B} \cdot d\vec{\ell} = 0 \Rightarrow B = 0 \quad 0 \leq d \leq A - a$$

$$\int_{\Gamma_3} \vec{B} \cdot d\vec{\ell} = \mu_0 (NI - NI) \Rightarrow B = 0 \quad d \geq A + a$$

$$\int_{\Gamma_1} \vec{B} \cdot d\vec{\ell} = \mu_0 NI = 2\pi r B \quad A - a < r < A + a$$

$$\Rightarrow B = \frac{\mu_0}{2\pi} \frac{NI}{d} = \frac{\mu_0}{2\pi} \frac{n 2\pi (A - a)}{d} I$$

$$\begin{array}{c} r \\ |d - A| < a \\ \uparrow \\ A - a < d < A + a \end{array}$$

$$B = \mu_0 n \frac{A - a}{A + r} I \quad \xrightarrow{A \rightarrow \infty} B = \mu_0 n I \quad \text{vně}$$

$$\Rightarrow B = \begin{cases} \mu_0 n I \vec{e}_z & r < a \\ 0 & r > a \end{cases}$$

→ na povrchu:

$$[\vec{n} \times \vec{B}] = \mu_0 I$$

$$\vec{n} = -\vec{e}_r \Rightarrow -\mu_0 n I \vec{e}_r \times \vec{e}_z = \mu_0 I \vec{e}_\varphi \Rightarrow \vec{l} = n I \vec{e}_\varphi = L \vec{e}_\varphi$$

$$\Rightarrow \vec{B} = \mu_0 L \vec{e}_z$$

→ Pre dva solenoidy s polomerami $a_2 < a_1$ bude pole:

$$\vec{B} = \begin{cases} \mu_0 (L_2 + L_1) & r < a_2 \\ \mu_0 L_1 & a_2 < r < a_1 \\ 0 & a_1 < r \end{cases}$$

10.3 Vektorový potenciál homo. mag. pole

$$\vec{B} = B \vec{e}_z \quad B_i = \epsilon_{ijk} \partial_j A_k$$

$$\vec{B} = \nabla \times \vec{A} \rightsquigarrow B = \partial_x A_y - \partial_y A_x = B, \quad A_z = 0$$

$$\begin{aligned} A_x &= By & A'_x &= -By \\ A_y &= 2Bx & A'_y &= 0 \end{aligned}$$

$$\vec{A} = B(y, 2x, 0) \quad \vec{A}' = B(-y, 0, 0)$$

$$\nabla \psi = B(2y, -2x, 0) \Rightarrow \psi = -2Bxy$$

$$\vec{A} = -\frac{1}{2} \vec{r} \times \vec{B} \Rightarrow A_i = -\frac{1}{2} \epsilon_{ijk} x_j B_k$$

$$\begin{aligned} \nabla \times \vec{A} &\rightarrow \epsilon_{ijk} \partial_j A_k = \epsilon_{ijk} \partial_j \left(-\frac{1}{2} \epsilon_{klm} x_l B_m \right) = -\frac{1}{2} \epsilon_{ijk} \epsilon_{klm} \delta_{lj} B_m = \\ &= \frac{1}{2} (\delta_j^l \delta_i^m - \delta_j^m \delta_i^l) \delta_{lj} B_m = \frac{1}{2} (\delta_j^j \delta_i^m - \delta_i^m) B_m = \\ &= \frac{1}{2} (3 - 1) B_i = B_i \end{aligned}$$

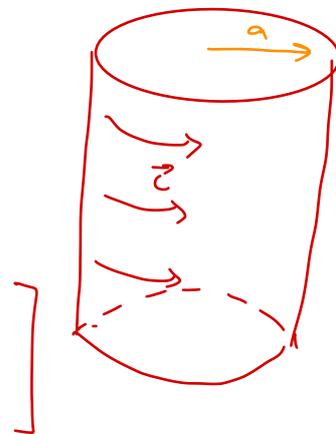
$$\psi = \int_S \vec{B} \cdot d\vec{S} = \int_S \nabla \times \vec{A} \cdot d\vec{S} = \oint_{\partial S} \vec{A} \cdot d\vec{l}$$

10.4 Vektorový potenciál nekonečného solenoidu

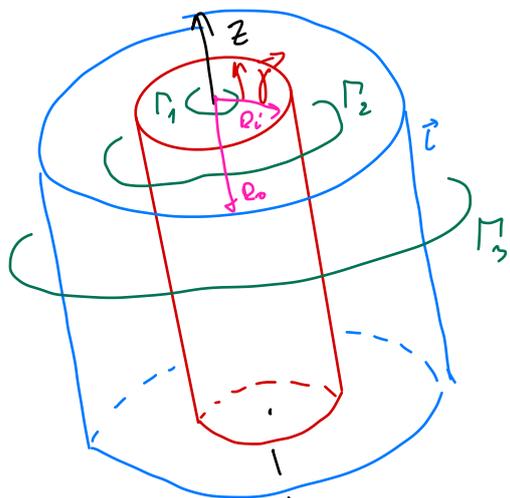
$$\begin{aligned} \vec{A} &= \frac{\mu_0}{4\pi} \int \frac{\vec{J}(\vec{r}')}{|\vec{r} - \vec{r}'|} dV = \\ &= \frac{\mu_0}{4\pi} \int_{-\infty}^{\infty} \int_0^{2\pi} \frac{c a}{|\vec{r} - \vec{r}'|} d\varphi dz \end{aligned}$$

$$|\vec{r} - \vec{r}'|^2 = (r \cos \varphi - a \cos \varphi')^2 + (r \sin \varphi - a \sin \varphi')^2 + (z - z')^2 = r^2 + a^2 - 2ra \cos(\varphi - \varphi') + (z - z')^2$$

$$= \frac{\mu_0 a c}{4\pi} \int_0^{2\pi} \int_{-\infty}^{\infty} \frac{1}{(r^2 + a^2 - 2ra \cos(\varphi - \varphi') + z'^2)^{\frac{1}{2}}} dz' d\varphi$$



10.5 Magnetické pole koaxiálního vodiče



$$\vec{j} = j\vec{e}_z \quad \vec{l} = i\vec{e}_z$$

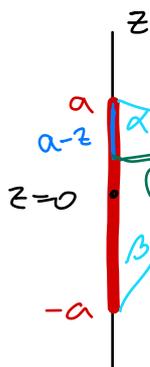
$$\oint_{\Gamma_3} \vec{B} \cdot d\vec{l} = \mu_0 (j \cdot S + I l) = 2\pi r B$$

$$\oint_{\Gamma_2} \vec{B} \cdot d\vec{l} = \mu_0 j S = 2\pi r B$$

$$\oint_{\Gamma_1} \vec{B} \cdot d\vec{l} = \mu_0 j \pi r^2 = 2\pi r B$$

$$\Rightarrow B_\varphi = \begin{cases} \frac{\mu_0}{2\pi} \frac{1}{r} (j\pi R_i^2 + 2\pi R_o i) = \frac{\mu_0}{2r} (j R_i^2 + 2R_o i) & r > R_i \\ \frac{\mu_0}{2\pi} \frac{1}{r} (j\pi R_i^2) = \frac{\mu_0}{2r} j R_i^2 & R_o < r < R_i \\ \frac{\mu_0}{2} j r & R_o > r \end{cases}$$

10.6 Magnetické pole v přímém úseku vodiče



$$\vec{B} = \frac{\mu_0 I}{4\pi} \int \frac{d\vec{l}' \times (\vec{x} - \vec{x}')}{|\vec{x} - \vec{x}'|^3}$$

$$\left. \begin{aligned} \vec{x}' &= (0, 0, z') \\ \vec{x} &= (r \cos \varphi, r \sin \varphi, z) \end{aligned} \right\} \vec{x} - \vec{x}' = \vec{e}_x r \cos \varphi + \vec{e}_y r \sin \varphi + (z - z') \vec{e}_z$$

$$d\vec{l}' = (0, 0, dz') = \vec{e}_z dz'$$

$$d\vec{l}' \times \vec{r}' = r \cos \varphi dz' \vec{e}_z \times \vec{e}_x + r \sin \varphi dz' \vec{e}_z \times \vec{e}_y$$

$$\vec{B} = \frac{\mu_0 I}{4\pi} \int_{-a}^a \frac{r \cos \varphi dz' \vec{e}_y - r \sin \varphi dz' \vec{e}_x}{(r^2 + (z - z')^2)^{3/2}}$$

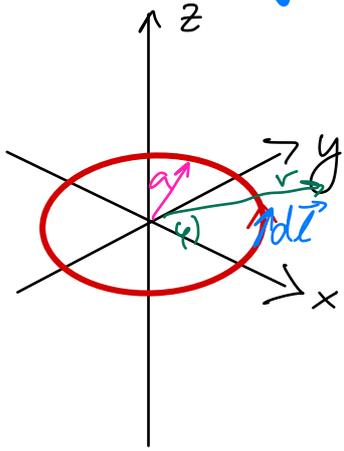
$$= \mu_0 \frac{I}{4\pi} (r \cos \varphi \vec{e}_y - r \sin \varphi \vec{e}_x) \int_{-a-z}^{a-z} \frac{dz'}{(r^2 + z'^2)^{3/2}}$$

$$= \mu_0 \frac{I}{4\pi} (r \cos \varphi \vec{e}_y - r \sin \varphi \vec{e}_x) \left[\frac{z'/R^2}{(R^2 + z'^2)^{1/2}} \right]_{R \cot \beta}^{R \cot \alpha} =$$

$$= \mu_0 \frac{I}{4\pi} (\cos \varphi \vec{e}_y - \sin \varphi \vec{e}_x) \left[\frac{R \cot \beta}{R(1 + \cot^2 \beta)^{1/2}} - \frac{R \cot \alpha}{R(1 + \cot^2 \alpha)^{1/2}} \right] = \mu_0 \frac{I}{4\pi R} (\cos \alpha - \cos \beta) \vec{e}_\varphi$$

$$\frac{\sin^2 + \cos^2}{\sin^2} = \frac{1}{\sin^2 \alpha}$$

10.7 Magnetické pole proudu v kruhové smyčce



$$d\vec{B} = \frac{\mu_0 I}{4\pi} \frac{d\vec{l} \times (\vec{r} - \vec{r}')}{|\vec{r} - \vec{r}'|^3}$$

$$\vec{r}' = (a \cos \varphi', a \sin \varphi', 0)$$

{ dotyčnice

$$d\vec{l}' = (-a \sin \varphi', a \cos \varphi', 0) d\varphi'$$

$$\vec{B} = \frac{\mu_0 I}{4\pi} \int_0^{2\pi} \frac{(-a \sin \varphi', a \cos \varphi', 0) \times (r \cos \varphi - a \cos \varphi', r \sin \varphi - a \sin \varphi', z) d\varphi'}{|\vec{r} - \vec{r}'|^3}$$

$$\begin{aligned} |\vec{r} - \vec{r}'|^2 &= (r \cos \varphi - a \cos \varphi')^2 + (r \sin \varphi - a \sin \varphi')^2 + (z - z')^2 = \\ &= r^2 + a^2 - 2ra \cos \varphi \cos \varphi' - 2ra \sin \varphi \sin \varphi' + (z - z')^2 = \\ &= r^2 + a^2 - 2ra \cos(\varphi - \varphi') + z^2 \end{aligned}$$

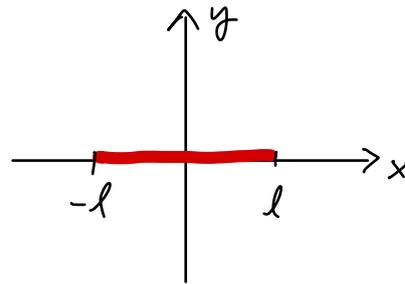
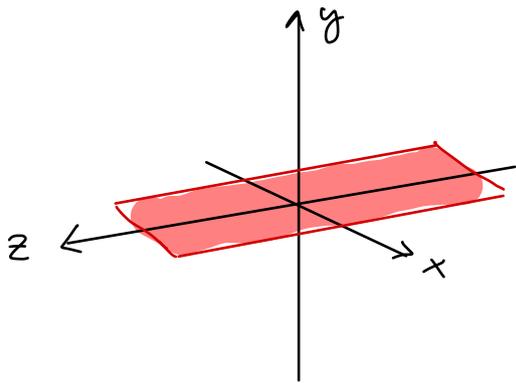
$$\begin{aligned} d\vec{l} \times (\vec{r} - \vec{r}') &= (a \cos \varphi' z, a \sin \varphi' z, \\ &\quad -a \sin \varphi' (r \sin \varphi - a \sin \varphi') - a \cos \varphi' (r \cos \varphi - a \cos \varphi')) = \\ &= (a \cos \varphi' z, -a \sin \varphi' z, -a r \cos(\varphi - \varphi') + a^2) \end{aligned}$$

$$\vec{B} = \frac{\mu_0 I}{4\pi} \int_0^{2\pi} \frac{(a z \cos \varphi', a z \sin \varphi' z, -a r \cos(\varphi - \varphi') + a^2)}{(r^2 + a^2 - 2ra \cos(\varphi - \varphi') + z^2)^{3/2}} d\varphi'$$

• na ose : $r = 0$

$$\begin{aligned} \vec{B} &= \frac{\mu_0 I}{4\pi} \int_0^{2\pi} \frac{(a z \cos \varphi', a z \sin \varphi' z, a^2)}{(a^2 + z^2)^{3/2}} d\varphi = \\ &= \frac{\mu_0 I}{2} \frac{(0, 0, a^2)}{(a^2 + z^2)^{3/2}} \end{aligned}$$

10.9 Magnetické pole proudu tekoucího plošným vodičem



$$\vec{c} = c(0, 0, 1)$$

$$\vec{B} = \frac{\mu_0 c}{4\pi} \int_{-\infty}^{\infty} \int_{-l}^l \frac{(0, 0, 1) \times (x-x', y, z-z')}{((x-x')^2 + y^2 + (z-z')^2)^{\frac{3}{2}}} dx' dz' =$$

$$= \frac{\mu_0 c}{4\pi} \int_{-\infty}^{\infty} \int_{-l}^l \frac{\vec{e}_y(x-x') - \vec{e}_x y}{((x-x')^2 + y^2 + (z-z')^2)^{\frac{3}{2}}} dx' dz' =$$

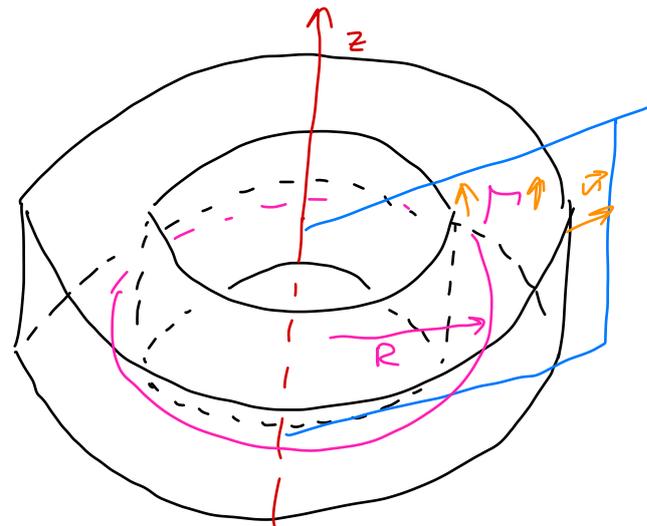
$$= \frac{\mu_0 c}{4\pi} \int_{-l}^l \vec{e}_y(x-x') - \vec{e}_x y \int_{-\infty}^{\infty} \frac{1}{((x-x')^2 + y^2 + (z-z')^2)^{\frac{3}{2}}} dz' dx'$$

$$\left[\frac{z' \sqrt{(x-x')^2 + y^2}}{((x-x')^2 + y^2 + z'^2)^{\frac{1}{2}}} \right]_{-\infty}^{\infty} = \frac{2}{(x-x')^2 + y^2}$$

$$= \frac{\mu_0 c}{2\pi} \int_{-l}^l \frac{\vec{e}_y(x-x') - \vec{e}_x y}{(x-x')^2 + y^2} dx' =$$

$$= \frac{\mu_0 c}{2\pi} \left(\vec{e}_y \frac{1}{2} \ln \left(\frac{(x+l)^2 + y^2}{(x-l)^2 + y^2} \right) - \vec{e}_x \left[\arctg \left(\frac{x-l}{y} \right) - \arctg \left(\frac{x+l}{y} \right) \right] \right)$$

10.10 Magnetické pole solenoidu



$$\vec{B} = B \vec{e}_\varphi$$

$$\oint \vec{B} \cdot d\vec{l} = I \mu_0$$

$$\int_{\Gamma} \vec{B} \cdot d\vec{l} = I \mu_0 \Rightarrow B = \frac{I \mu_0}{2\pi R}$$

$$\oint \vec{B}_{\text{vni}} \cdot d\vec{l} = 0 \Rightarrow \vec{B}_{\text{vni}} = 0$$

$$\vec{B} = \frac{I \mu_0}{2\pi R} \vec{e}_\varphi \quad \partial_\varphi B = 0 \quad \nabla \times \vec{B} \Rightarrow \frac{1}{R} \partial_R (RB) = 0 \quad -\partial_z B = 0$$

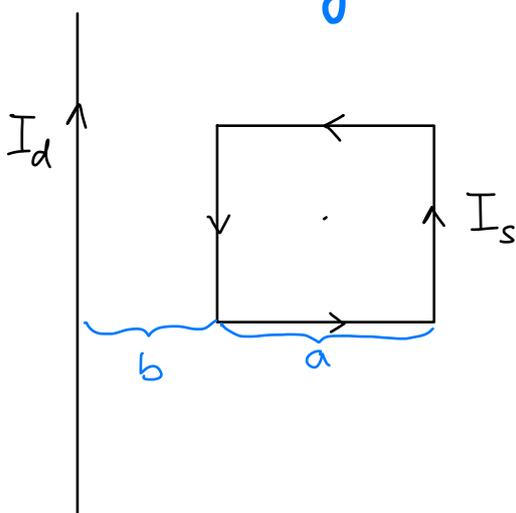
$$\vec{n} = n_R \vec{e}_R + n_z \vec{e}_z$$

$$\vec{n} \times \vec{B} = \frac{I \mu_0}{2\pi R} \vec{n} \times \vec{e}_\varphi = \mu_0 \vec{C} \Rightarrow \vec{C} = \frac{I}{2\pi R} \vec{n} \times \vec{e}_\varphi = C \vec{n} \times \vec{e}_\varphi =$$

$$= \frac{I}{2\pi R} (n_R \vec{e}_z - n_z \vec{e}_R)$$

$$\vec{n} \cdot \vec{B} = (\dots) \vec{e}_R \cdot \vec{e}_\varphi + (\dots) \vec{e}_z \cdot \vec{e}_\varphi = 0$$

10.11 Přímý vodič a čtvercová proudová smyčka



$$\vec{B}_d = \frac{\mu_0 I_d}{2\pi} \frac{1}{r} \vec{e}_\varphi = B \vec{e}_\varphi$$

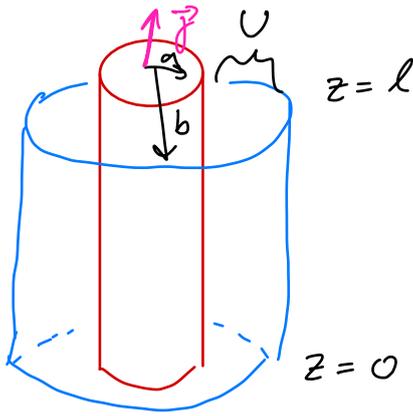
to symetrie

$$\vec{F} = I_s \int d\vec{S} \times \vec{B} = I_s \int_b^{b+a} d\vec{r} \times \vec{B} + I_s \int_{b+a}^b d\vec{r} \times \vec{B}$$

$$+ I_s \int_0^a B(b+a) \vec{e}_z \times \vec{e}_\varphi dz + I_s \int_a^0 B(b) \vec{e}_z \times \vec{e}_\varphi dz =$$

$$= -I_s \vec{e}_r \frac{\mu_0 I_d a}{\pi} \left(\frac{1}{b+a} - \frac{1}{b} \right)$$

10.13 Koaxiální vodič zapojený do obvodu



$$\vec{E} = 0 \quad r > b$$

$$r\vec{E} = \vec{j} = j\vec{e}_z \quad r < a$$

$$a < r < b : \frac{1}{r} \partial_r (r \partial_r \phi) + \frac{1}{r^2} \frac{\partial^2 \phi}{\partial \varphi^2} + \frac{\partial^2 \phi}{\partial z^2} = 0$$

$$\Rightarrow r \partial_r \phi = c \Rightarrow \phi = c \ln \frac{r}{r_0} \Rightarrow \phi = cz \ln \frac{r}{r_0}$$

$$\phi(z=l, r=b) = c l \ln \frac{b}{r_0} = 0 \Rightarrow r_0 = b$$

$$\phi(z=l, r=a) = c l \ln \frac{a}{b} = U \Rightarrow c = \frac{U}{l \ln \frac{a}{b}} \Rightarrow \phi = \frac{Uz}{l} \frac{\ln \frac{r}{b}}{\ln \frac{a}{b}}$$

$$\vec{E} = -\nabla \phi = -\frac{U}{l} \frac{\ln \frac{r}{b}}{\ln \frac{a}{b}} \vec{e}_z - \frac{Uz}{l r \ln \frac{a}{b}} \vec{e}_r \quad a < r < b$$

$$[\vec{n} \times \vec{E}] = 0 \Rightarrow \frac{j}{r} - \frac{U}{l} = 0 \Rightarrow j = \frac{Ur}{l} \Rightarrow I = \frac{U}{l} \pi a^2$$

$$\Rightarrow R = \frac{l}{j \pi a^2}$$

11.1 Magnetické pole dipólu

$$\vec{A} = \frac{\mu_0}{4\pi} \frac{\vec{m} \times \vec{r}}{r^3} \quad \rightarrow \quad A_i = \frac{\mu_0}{4\pi} \frac{1}{r^3} \varepsilon_{ijk} m_j x_k \quad \frac{1}{(\dots)^{3/2}}$$

$$B_i = \varepsilon_{ijk} \partial_j A_k = \frac{\mu_0}{4\pi} m_l \varepsilon_{ijk} \varepsilon_{klm} \partial_j \frac{x_m}{r^3} =$$

$$= \frac{\mu_0}{4\pi} m_l (\delta_i^l \delta_j^m - \delta_i^m \delta_j^l) \frac{\delta_j^m r^3 - x_m 3r x_j}{r^6} =$$

$$= \frac{\mu_0}{4\pi} (\delta_j^m m_i - \delta_i^m m_j) \frac{\delta_j^m r^2 - 3 x_m x_j}{r^5} =$$

$$= \frac{\mu_0}{4\pi} \left(\frac{3r^2 m_i - 3r^2 m_i}{r^4} - \frac{m_i r^2 - 3 x_i x_j m_j}{r^5} \right) =$$

$$= \frac{\mu_0}{4\pi} \frac{1}{r^3} \left(\frac{3 x_i x_j m_j}{r^2} - m_i \right) \quad \rightarrow \quad \vec{B} = \frac{\mu_0}{4\pi} \frac{1}{r^3} \left(\frac{3 \vec{r} (\vec{r} \cdot \vec{m})}{r^2} - \vec{m} \right)$$

11.2 Dipól-dipól interakcia - sila

$$\vec{B}_1|_2 = \frac{\mu_0}{4\pi} \frac{1}{r_{12}^3} \left(\frac{3 \vec{r}_{12} (\vec{r}_{12} \cdot \vec{m}_1)}{r_{12}^2} - \vec{m}_1 \right)$$

$$\vec{B}_1 \cdot \vec{m}_2 = \frac{\mu_0}{4\pi} \frac{1}{r_{12}^3} \left(3 \frac{(\vec{r}_{12} \cdot \vec{m}_2)(\vec{r}_{12} \cdot \vec{m}_1)}{r_{12}^2} - \vec{m}_1 \cdot \vec{m}_2 \right) = \frac{\mu_0}{4\pi} \frac{1}{r_{12}^3} \mathcal{M}$$

$$\vec{F} = \nabla (\vec{B} \cdot \vec{m}_2) = \frac{\mu_0}{4\pi} \left(\mathcal{M} \nabla \frac{1}{r_{12}^3} + \frac{1}{r_{12}^3} \nabla \mathcal{M} \right)$$

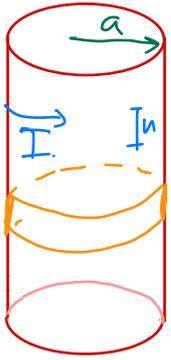
$$-3 \frac{\vec{r}}{r_{12}^5} \quad -2 \frac{\vec{r}}{r_{12}^4}$$

$$\nabla \mathcal{M} = \partial_i \left(3 \frac{x_j m_j^2 x_l m_l^1}{r^2} - m_1^i m_2^k \right) = 3 (\vec{r}_{12} \cdot \vec{m}_2) (\vec{r}_{12} \cdot \vec{m}_1) \nabla \frac{1}{r^2} + \frac{3}{r^2} m_j^2 m_l^1 (\delta_{ij} x_l + \delta_{il} x_j) =$$

$$= -6 (\vec{r}_{12} \cdot \vec{m}_2) (\vec{r}_{12} \cdot \vec{m}_1) \frac{\vec{r}}{r^4} + \frac{3}{r^2} \left[\vec{m}_2 (\vec{m}_1 \cdot \vec{r}_{12}) + \vec{m}_1 (\vec{m}_2 \cdot \vec{r}_{12}) \right]$$

$$\vec{F} = \frac{\mu_0}{4\pi} \left(-15 (\vec{r}_{12} \cdot \vec{m}_2) (\vec{r}_{12} \cdot \vec{m}_1) \frac{\vec{r}}{r^7} + \frac{3}{r^5} \left[\vec{m}_2 (\vec{m}_1 \cdot \vec{r}_{12}) + \vec{m}_1 (\vec{m}_2 \cdot \vec{r}_{12}) \right] + 3 \frac{\vec{r}}{r^3} \vec{m}_1 \cdot \vec{m}_2 \right)$$

12.1 Samoindukčnost solenoidu



$$LB = \mu_0 NI \Rightarrow B = \mu_0 n I$$

$$W_M = \frac{1}{2\mu_0} \int B^2 dV = \frac{\mu_0}{2} n^2 I^2 \int dV = \frac{\mu_0}{2} n^2 I^2 \pi a^2 d$$

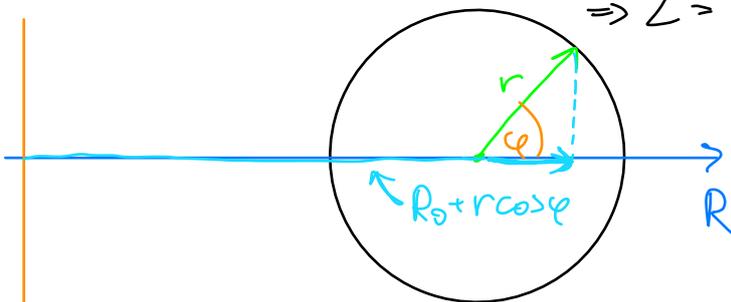
$$\Rightarrow L = \mu_0 n^2 \pi a^2 d \Rightarrow \ell = \mu_0 n^2 a^2 \pi$$

12.2 Samoindukčnost toroidalnika solenoidu

$$\vec{B} = \frac{I \mu_0}{2\pi R} \vec{e}_\varphi$$

$$W = \frac{1}{2\mu_0} \int B^2 dV = \frac{\mu_0 I^2}{8\pi^2} \iint \frac{1}{R^2} dl dS = \frac{\mu_0 I^2}{8\pi^2} \int \frac{2\pi R}{R^2} dS$$

$$\Rightarrow L = \frac{\mu_0}{2\pi} \int \frac{1}{R} dS$$



a) kruh

$$L_{\text{kruh}} = \frac{\mu_0}{2\pi} \int \frac{1}{R} dS = \frac{\mu_0}{2\pi} \int_0^{a_0} \int_0^{2\pi} \frac{r}{R_0 + r \cos \varphi} dr d\varphi = \begin{cases} u = R_0 + r \cos \varphi \\ du = dr \cos \varphi \end{cases}$$

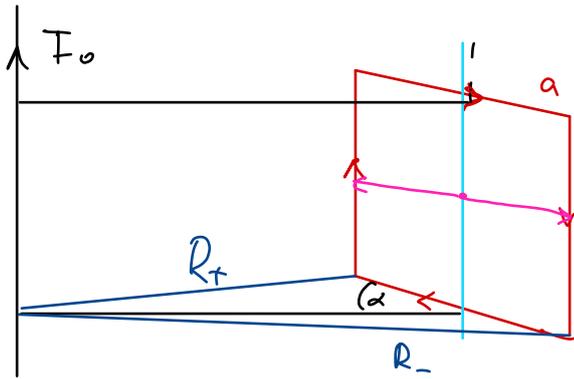
$$= \frac{\mu_0}{2\pi} \int_0^{2\pi} \int_0^{a_0} \frac{u - R_0}{\cos \varphi} \frac{du}{\cos \varphi} d\varphi =$$

$$= \frac{\mu_0}{2\pi} \int_0^{2\pi} \frac{1}{\cos^2 \varphi} \int_{R_0}^{R_0 + a_0 \cos \varphi} \left(1 - \frac{R_0}{u}\right) du d\varphi = \frac{\mu_0}{2\pi} \int_0^{2\pi} \frac{1}{\cos^2 \varphi} \left(a_0 \cos \varphi - R_0 \ln \left(1 + \frac{a_0 \cos \varphi}{R_0}\right) \right) d\varphi$$

c) obdelnik

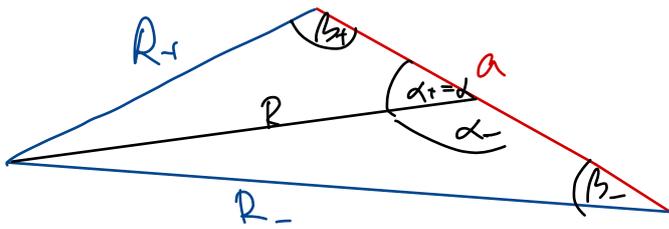
$$L_{\text{obdelnik}} = \frac{\mu_0}{2\pi} \int_a^b \int_0^d \frac{1}{R} dR dz = \frac{\mu_0 d}{2\pi} \ln \frac{b}{a}$$

Moment síly působící na obdelníkový smyčku



$$\vec{B} = \frac{\mu_0}{2\pi} \frac{I_0}{R} \vec{e}_\varphi$$

$$d\vec{s} = (0, 0, dz)$$



$$\frac{\sin \mu_{\pm}}{R} = \frac{\sin \alpha_{\pm}}{R_{\pm}} \rightarrow \sin \mu_{\pm} = \frac{\sin \alpha_{\pm} R}{R_{\pm}}$$

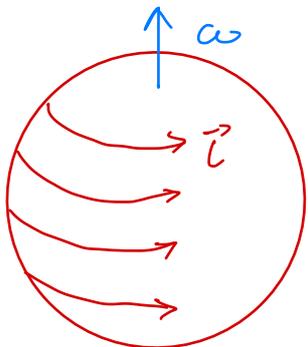
$$\vec{F} = I \int d\vec{s} \times \vec{B} = \frac{\mu_0 I_0 I}{2\pi R_{\pm}} \int_0^b \vec{e}_\varphi \times \vec{e}_z = \mp \frac{\mu_0 I_0 I b}{2\pi R_{\pm}} \vec{e}_R$$

$$\vec{M}_{\pm} = \vec{d}_{\pm} \times \vec{F}_{\pm} = \frac{a}{2} F_{\pm} \sin \mu_{\pm} \vec{e}_z = \mp \frac{a}{2} \frac{\mu_0 I_0 I b}{2\pi} \frac{\sin \alpha R}{R_{\pm}^2}$$

$$R_{\pm}^2 = R^2 + \frac{a^2}{4} - R a \cos \alpha_{\pm} = R^2 + \frac{a^2}{4} \mp R a \cos \alpha$$

$$\vec{M} = \vec{M}_+ + \vec{M}_- = \frac{a}{2} \frac{\mu_0}{2\pi} a b I_0 I \sin \alpha \left(\frac{R}{R_+^2} + \frac{R}{R_-^2} \right) \vec{e}_z$$

Rotující nabitá koule



$$I = \sigma V = \sigma \omega R = \frac{Q}{4\pi R^2} R \omega = \frac{Q\omega}{4\pi R}$$

$$\vec{m} = \frac{1}{2} \int \vec{r} \times \vec{j} dV = \frac{1}{2} \int \vec{r} \times \vec{I} dS$$

$$\vec{r} = \vec{e}_r + r \sin \vartheta \vec{e}_\varphi + r \vec{e}_{\vartheta}$$

$$\vec{I} = \frac{Q\omega}{4\pi R} \vec{e}_\varphi$$

